Add Hole Information Loss Paradoll and the AnaBHEL Experiment

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Tribute to Stephen Hawking (1942-2018)



Black hole Hawking evaporation – Connecting GR, QM, SM in one stroke





Information Loss



- Entanglement of Hawking radiation?
- Firewall?
- Fussball?
- Etc., etc., etc.





End point

48 years after Hawking's discovery, the nature of BH evaporation is still under debate!

Can Hawking radiation carry information out after all? If so, how?



Hawking evaporation may result in the loss of

- First pointed out by Hawking himself in 1976
- Endless debates ever since
- Solutions include "black hole complementarity" (Susskind et al.), Firewalls (AMPS, AMPSS 2012), etc.
- Entanglement between Hawking radiation and partner particles (Wilczek 1987, Schutzhold-Unruh 2010, Hotta-Schutzhold-Unruh 2015)
- Planck size black hole remnants (Chen-Ong-Yeom, Phys. Rep. 2015)
- Naked black hole firewalls (Chen-Ong-Page-Sasaki-Yeom, PRL 2016)
- BMS Soft Hairs (Hawking-Perry-Strominger, 2016)
- No firewalls & nothing wrong w. information loss (Unruh-Wald 2017)
- Islands and replica wormholes based on string theory (Almheiri et al. 2019)
- Wave function of the universe and multi-histories (Chen-Sasaki-Yeom-Yoon 2021)









"Quantum entanglement is not just a property of QM, it is THE character of QM. It fundamentally breaks QM from classical physics." (E. Schrodedinger) ⁷

Monogamy of quantum entanglement



When would BH information come out? $S(B) = \log m$ (For a pure and random system, $S(B|A) = \sum_{n=1}^{m} \frac{1}{n} - \frac{m-1}{n} \text{ conjectured by Page, 1993; proved}$ by Sen, 1996.) 2nIf $S(A) \propto Area$, then the $S(A \cup B) = \log N = const$ information will come S(B)out when the black hole Informati Youn Old initial area decreases to $I = S(B) - S(B \mid A)$ Black g half value. This is called Black Høle th Hole

S(A)



Pure state black hole Mixed state

Page/Time

Resolution with Euclidean path integral

(P. Chen, M. Sasaki D -h. Yeom I. Yoon arXiv:2111.01005, 2021)



FIG. 1: Left: The Page curve, where S_A and S_B are Boltzmann entropies of A and B, respectively, and $S_B(A)$ denotes the entanglement entropy. Right: The path integral from the in-state $|in\rangle$ to the out-state $|out\rangle$, where the out-state is a superposition of classical boundaries $\{|i\rangle\}$.

Multi-histories and late-time dominance conditions of BH evolution



FIG. 2: Left: the causal structure of the usual semi-classical black hole, where the green curve is the trajectory of the collapsing matter, the red curve is the apparent horizon, and the blue line is the event horizon. Middle: the causal structure after a quantum tunneling at the time slice t. After the tunneling, matter or information (red curve) is emitted and the black hole structure disappears. Right: h_1 is the information-losing history, while $h_2^{(1,2)}$ are the information-preserving histories. Tunneling may happen either early time (A) or late time (B); the tunneling probability must be dominated at the late time.

Page curve modified



FIG. 3: Left: Probabilities of the information-losing history $(p_1, \text{ blue curve})$ and the informationpreserving history $(p_2, \text{ black curve})$. Right: Entropy of emitted radiation S_{rad}/S_0 vs. entanglement entropy S_{ent}/S_0 , where $S_0 = 3$ (black), 10 (blue), 50 (red), respectively. The thin red dashed curve is the location of the Page time, i.e., $dS_{\text{ent}}/dS_{\text{rad}} = 0$.

Comparison with recent developments based on string theory

- The string-based islands approach and ours share two essential features.
 - There exist more than two contributions to the evolution of the entanglement entropy, where one is dominant at early-time and the other at late-time.
 - 2. The final entanglement entropy is dominated by the late-time condition.

Analogous in spirit to the Quantum Extremal Surface (QES) computations for BHs.

 However, the entanglement entropy computed in QES is based on the density matrix instead of the quantum states, which is the standard quantum field theory approach and what we have followed.

theoretical Astro black holes too cold and too

Lifetime of solar BH: 10⁶⁷ years Age of the universe: 1.38 x 10¹⁰ years



Analog Black Holes

- Sound waves in moving fluids "dumb holes" Unruh (1981, 1995)
- Traveling index of refraction in media Yablonovitch (1989)
- Violent acceleration of electron by lasers Chen-Tajima (1999)
- Electromagnetic waveguides Schutzhold-Unruh (2005)
- Bose-Einstein condensate
 Steinhauer (2014)
- Accelerating mirror Fulling-Davies (1976), Davies-Fulling-Unruh (1977), Wilczek (1989), Carlitz-Willey (1987), Hotta-Schutzhold-Unruh (2015), Chen-Mourou (2017), Chen-Yeom (2017)

Testing thermal nature of Hawking radiation

Can also Test ILP

Flying Mirror: Entanglement between Hawking &



Red-shift by a mirror



- One can calculate the **out-going energy flux** as a function of the mirror trajectory (for 2D spacetime).
- Definitely, there should be no information loss in the mirror dynamics.
- Then what can we learn about the **entanglement** 17

Mirror trajectories

Chen, D-h Yeom, "Entropy evolution of moving mirrors and the information loss probl RD 96, 025016 (2017) [arXiv:1704.08613]

• Using this formula, we can test several candidate traject $S(u) = -\frac{c}{12} \log \dot{p}(u)$



$$\begin{aligned} 12^{10\,\mathrm{g}\,\mathrm{P}(\mathrm{u})} & 12^{10\,\mathrm{g}\,\mathrm{P}(\mathrm{u})} \\ \frac{S(t)}{dt} &= A \sin^2 \pi \frac{t}{t_{\mathrm{P}}} & 0 \le t < t_{\mathrm{P}}, \\ &= -A \frac{t_{\mathrm{P}}}{t_{\mathrm{f}} - t_{\mathrm{P}}} \sin^2 \pi \frac{t - t_{\mathrm{P}}}{t_{\mathrm{f}} - t_{\mathrm{P}}} & t_{\mathrm{P}} \le t < t_{\mathrm{f}}, \end{aligned}$$

– Suddenly stopping mirror: $t_f = 15$,

– Slowly stopping mirror: $t_{\rm f} = 20$,

– Long propagating mirror: $t_{\rm f} = 50$.

Q: How to physically realize it? A: Plasma Wakefield Acceleration

Tajima-Dawson (1979)- Laser driven (LWFA)

Chen-Dawson-Huff-Katsouleas (1985)- Particle beam driven (PWFA)



- SLAC & LBL- Acceleration of O(100) GeV/m observed!
- AWAKE- A new PWFA experiment at CERN

By now plasma wakefield acceleration has become a very active field of research worldwide.

Plasma wakefield is like a tsunami



Relativistic Flying Plasma Mirror

Bulanov (2001), Bulanov, Esirkepov, Tajima (2003), Mourou-Tajima-Bulanov (2006)

Reflected laser pulse Lorentz-boosted and tighter-focused.



Accelerating plasma mirror as an analog black P. Chen and G. Mourol, P.S. Rev. Lett. (2017).

SIMULATING A BLACK HOLE ON A TABLE

New black hole simulator may shed more light on a contradiction in fundamental physics



Black hole Hawking evaporation



Plasma wakefield speeds up as plasma density gradually decreases

Plasma density-Mirror trajectory Correspondence

P. Chen and G. Mourou. Phys. Plasmas (2020)

• Constant-plus-Gaussian density profile

 $n_p(x) = n_{p0}(a + be^{-x^2/2D^2})^2, \quad 0 \le x \le L,$

asymptotically gives rise to the Davies-Fulling trajectory:

 $x_M(t) = -\eta_a ct - Ae^{-\eta_a ct/D} + A, \quad t \to \infty, \quad \eta_a = 1 - a^2 \omega_{p0}^2 / 2\omega_0^2$ which in turn give the analog Hawking temperature

$$k_B T_H = \frac{\hbar c}{4\pi} \frac{\eta_a}{D} \,.$$

 Diff. BH scenarios corresponds to diff. trajectories, which can be tested thru pre-designed density profiles.

AnaBHEL Collaboration (Analog Black Hole Evaporation via Lasers)

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AnaBHEL Collaboration (Analog Black Hole Evaporation via Lasers) National Taiwan University + Ecole Polytechnique + Kansai Photon Science Inst.

• Three stages:

1. R&D of key components at NTU, Taiwan. Major

progress

2. KPSI laser facility @PW. Scheduled for summer of 2022

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Schematic Layout of AnaBHEL Experiment Entanglement between Hawking and partner photons A flying Einstein-Podolsky-Rosen experiment



Status of AnaBHEL R&D

- 1. Laser-plasma interaction dynamics
- 2. Deviation plasma mirrors & estimation of Hawking photon yield
- 3. Supersonic gas jet with density gradient
- 4. Superconducting nanowire single-photon Hawking detector

Trajectory of the flying plasma mirror with graded density

P. Chen and G. Mourou. Phys. Plasmas (2020)



- Flying plasma mirror is always one "bubble length" behind the driving laser.
- Bubble width can be controlled by tailoring the density of the plasma background, therefore the trajectory of the mirror can be designed.

Reflectivity of flying plasma mirror

Y.K. Liu, P. Chen and P. Fang. Phys. Plasmas (2021)

- Reflectivity *highly* depends on the density distribution of electrons that compose the flying plasma mirror.
- Square Root Lorentzian
 Distribution (SRLD) provides
 a better estimation of
 reflectivity.



Reflectivity of the relativistic flying mirror as a function of the source pulse frequency .

Semitransparency and finite-size effect of flying plasma mirror

Number of particles per mode for a semitransparent mirror in (1+3)D with any prescribed trajectory: $\frac{dN}{d^3k} \approx \int d^3p \left| \frac{\langle \mathbf{k}, \mathbf{p}; out | 0; in \rangle}{\langle 0; out | 0; in \rangle} \right|^2, \quad \beta_{\mathbf{k}\mathbf{p}} = (u_{\mathbf{p}}^{in^*}, u_{\mathbf{k}}^{out})_{KG} \approx \frac{\langle \mathbf{k}, \mathbf{p}; out | 0; in \rangle}{\langle 0; out | 0; in \rangle}.$

The result for a Dirac delta

mirror of infinite transverse area.

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Mimic the mirror by: $V(x) = \gamma^{-1}(t)H(\mathbf{x}_{\perp})f(z - Z(t))$

- Mirror's transversal distribution: $H(\mathbf{x}_{\perp})$ •
- Mirror's longitudinal distribution: f(z Z(t))

• Mirror's longitudinal distribution:
$$f(z - Z(t))$$

• Mirror's longitudinal distribution: $f(z - Z(t))$
We then have: $\beta_{\mathbf{kp}} \approx F(\mathbf{k}, \mathbf{p}) \times \left(\frac{-i\alpha}{16\pi^3 \sqrt{\omega_k \omega_p}} \int dt \gamma^{-1}(t) e^{i(\omega_k + \omega_p)t - i(k_3 - p_3)Z(t)}\right)$,

where $F(\mathbf{k}, \mathbf{p}) = \left| d^2 x_{\perp} H(\mathbf{x}_{\perp}) e^{-i(k_{\perp}+p_{\perp})\cdot x_{\perp}} \right| d\zeta f(\zeta) e^{-i(k_3-p_3)\zeta}, \quad \zeta = z - Z(t).$

Transverse distribution leads to diffraction, whereas longitudinal distribution may lead to enhancement of

Estimation of Hawking Photon Yield

The analog Hawking particles emitted by this semitransparent mirror has the particle spectrum [33] [gr-qc].

$$\frac{dN_{ref}}{d\omega_k d\Omega} \approx \frac{\alpha^2 \mathcal{F}_L(\mathbf{k}_\perp) \mathcal{F}_D(k_3)}{8\pi\kappa} \left[\frac{\omega_k}{e^{\omega_k/T_{eff}(\theta_k)} + 1} \right], \ T_{eff}(\theta_k) = \frac{\kappa}{(1 + \cos\theta_k)\pi}$$

where *F*'s are form factors due to the mirror's finite area & thickness, which are independent of reflectivity and mirror's

• Motiver of analog Hawking particles emission per laser

(
$$N \approx \int_{0}^{\kappa} d\omega_{k} \int d\Omega \frac{dN}{d\omega_{k} d\Omega} = (0.27 + 0.02),$$

Assuming 1 laser shot/min and 8-
hr running time/day, a 20-day
experiment with an ideal detector
efficiency gives the detected event
 $N_{detect} = (1 \times 60 \times 8 \times 20) \times 1 \times N \approx 3000.$

Significant Progress in R&D at NTU

Supersonic gas jet



Delivers expected density gradient





Superconducting nanowire singlephoton Hawking detector Near 100% quantum efficiency attainable!



Summary

- Hawking evaporation and information loss paradox is one of the fundamental problems in physics.
- So far investigations are essentially theoretical; Direct observation of black hole end-stage unlikely.
- Quantum entanglement between Hawking radiation and partner particle may reveal the secrete.
- Accelerating plasma mirrors may serve to address some aspects of this paradox experimentally.
- AnaBHEL experiment is in good progress.