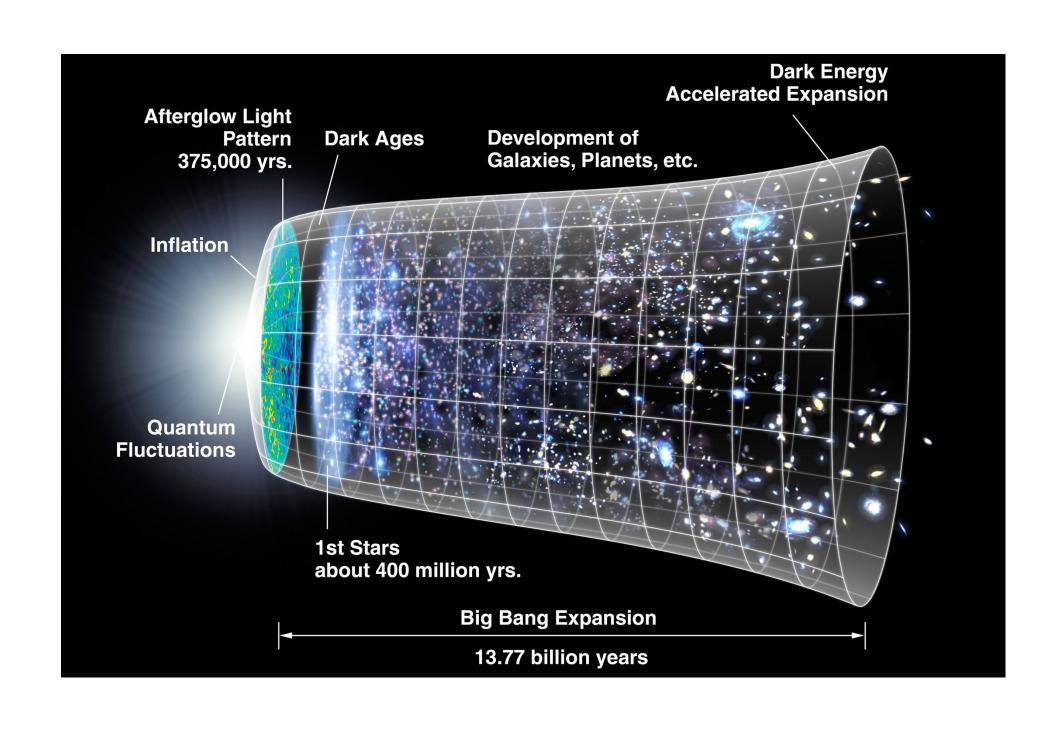
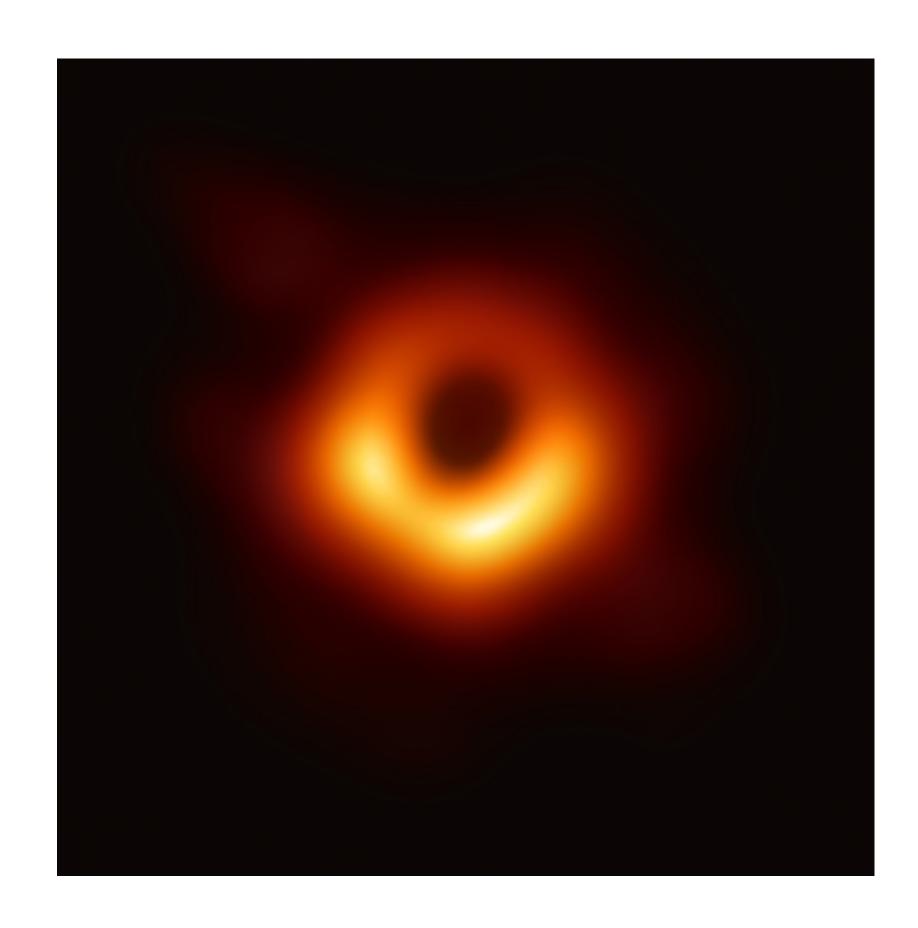
Renormalisation of essential operators in Asymptotic Safety - a status report

Benjamin Knorr







lack of smoking gun quantum gravity experiments

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many big ideas and even bigger claims on QG

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why trust any approach in particular?

recipe for a falsifiable and predictive quantum gravity theory:

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 - 1. set up quantum theory of gravity and matter (at least SM)
 - 2. confront the theory with as much available theory constraints (unitarity, causality, ...) and experimental data (cosmological evolution, particle masses, ...) as possible
 - 3. if consistent, only then move on to the "big questions": black holes, big bang, ...

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- less glamorous, more down-to-earth: fix starting point and see how far you get

- working hypothesis: QFT works also for QG
 - access to standard QFT notions like renormalisation group and scattering amplitudes
- in this talk: work towards $(2 \rightarrow 2)$ amplitudes in Asymptotic Safety, but general ideas can be carried over to other approaches

Outline

- Scattering amplitudes, field redefinitions and (in-)essential couplings
- Essential couplings in quantum gravity
- RG of essential QG selected results
- Summary

Scattering amplitudes, field redefinitions and (in-)essential couplings

- within Asymptotic Safety: most investigations at level of off-shell RG flow
 - residual dependencies on parameterisation, gauge choice etc.
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Effective Action

- for now: talk about **exact** theory (later: approximations) in the "not too strongly" interacting regime
 - hierarchy in the importance of correlation functions
 - on-shell expansion central to make sense of this
- recall: effective action includes all quantum effects, full scattering amplitudes obtained from tree level diagrams

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- consider "standard" scalar field theory in Minkowski space
 - wave function renormalisation inessential: $\phi \mapsto Z(p^2)^{-1/2}\phi$
 - removes non-trivial momentum dependence of propagator by shifting it into interactions
 - caveat: needs Z > 0, cannot remove or add modes
 - what about branch cuts (e.g. 1-loop logs)?

infinite derivative scalar field theory in Minkowski space:

$$\Gamma = \int d^d x \, \frac{1}{2} \phi \frac{\Box^2 + (m_1^2 + m_2^2)\Box + m_1^2 m_2^2}{2\Box + m_1^2 + m_2^2} \phi + \mathcal{O}(\phi^4)$$

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⇒ two non-ghostly modes

- consider Starobinsky-type theory
 - spectrum: EH + massive scalar
 - field redefinition can change value of \mathbb{R}^2 coupling, but cannot be set to zero (removes mode)

- in practice: have to choose universality class that we are investigating beforehand
 - adapt regularisation, make specific choices for inessential couplings to simplify computations
 - hierarchy expansion about on-shell configuration
 - how to treat branch cuts, (strong) IR non-localities, bound states, ...?

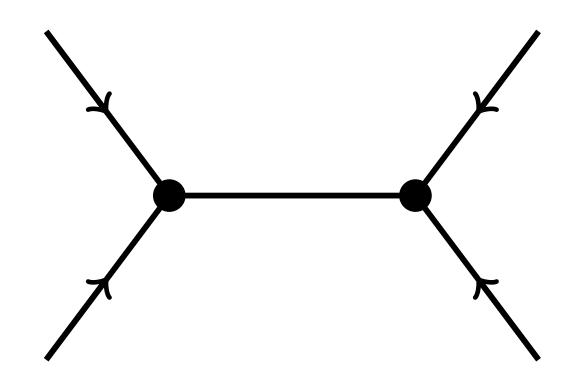
- in practice: have to choose universality class that we are investigating beforehand
- interpretation of approximation:
 - choose UC → adapt field redefinition to remove inessential operators → approximate
 - conceptually different from first approximating (introduces extra modes), then field redefinition (potentially removing extra modes)

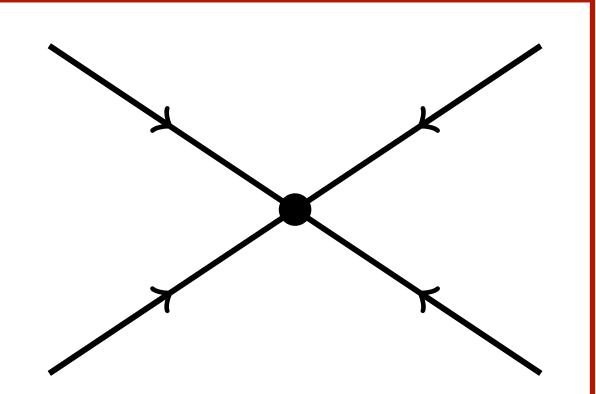
especially relevant in derivative expansions: Platania, Wetterich '20

Platania '22

- in practice: have to choose universality class that we are investigating beforehand
- interpretation of approximation
- in the following: discuss [GR] (+matter in [Gauss])
 - only "standard" modes (no extra modes, ghostly or not)

- in practice: have to **choose universality class** that we are investigating beforehand
- interpretation of approximation
- in the following: discuss [GR] (+matter in [Gauss])
- philosophy: any inessential couplings that can be set to zero should be set to zero ("minimal essential scheme")





Essential couplings in quantum gravity

Essential couplings in [GR]

general field redefinition of the metric:

$$g_{\mu\nu} \mapsto c_0 g_{\mu\nu} + c_1(\Box) R_{\mu\nu} + c_2(\Box) R g_{\mu\nu} + \dots$$

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 (is this still a Lorentzian metric?)

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• curvature expansion of effective action of [GR] in minimal scheme:

$$\Gamma_{[\mathbf{GR}]} = \frac{1}{16\pi G_N} \int d^d x \left[2\Lambda - R + \Theta \mathfrak{E} + G_{C^3} C^3 + \mathcal{O}(\mathcal{R}^4) \right]$$

NB: only one of G_N , Λ essential, cannot set either to zero generically

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quadratic terms inessential: multiply EoM

$$\left(R - \frac{2d}{d-2}\Lambda + \dots\right) f_R(\square) \left(R - \frac{2d}{d-2}\Lambda + \dots\right)$$
$$\left(S_{\mu\nu} + \dots\right) f_S(\square) \left(S^{\mu\nu} + \dots\right)$$
$$\square C = \{DDS, DDR\}$$

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- quadratic terms inessential: multiply EoM
- cubic terms almost all inessential, except Goroff-Sagnotti term

no essential $D^2C^3, ..., D^{12}C^3$ terms (Bianchi + partial integration)

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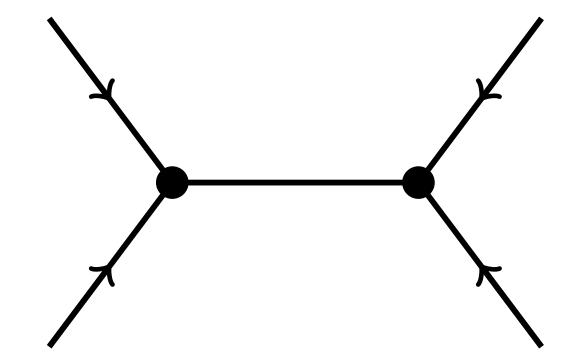
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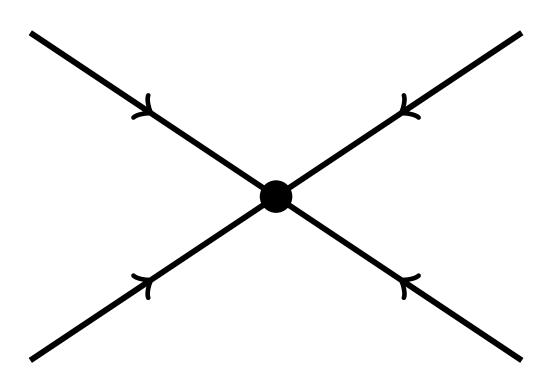
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- essential higher order terms: **only** Weyl + covariant derivatives, no $\ \square\ C$

structurally,
$$f(s, t) C^4$$

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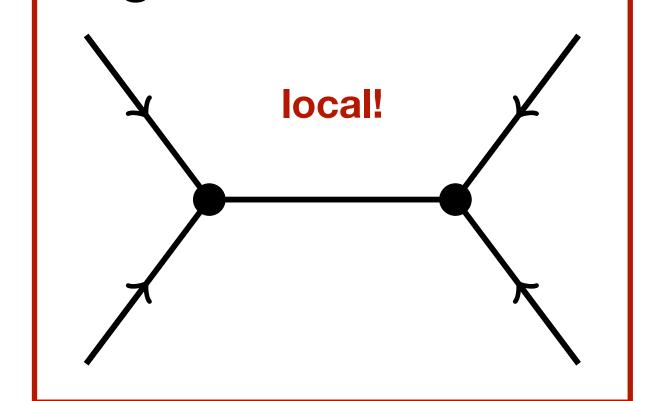


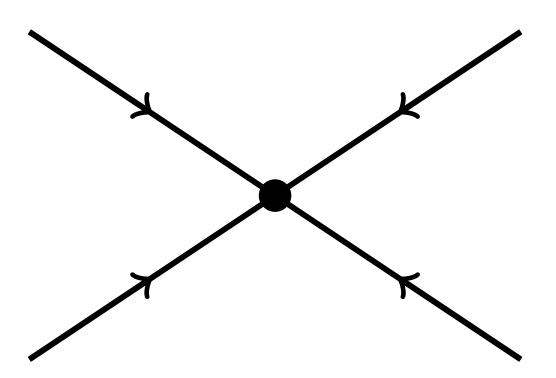


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- if you are in [GR]: **non-trivial** scattering physics starts at four-point function (big oof)
- in practice: have to impose field redefinition consistently along RG flow reduction of complexity, but no free lunch
- non-minimal coupling to matter: only involves Weyl tensor → expected to be extremely suppressed

$$C^{\mu\nu\rho\sigma}F_{\mu\nu}F_{\rho\sigma}$$

RG of essential QG selected results

• consider derivative expansion up to 6th order (d=4):

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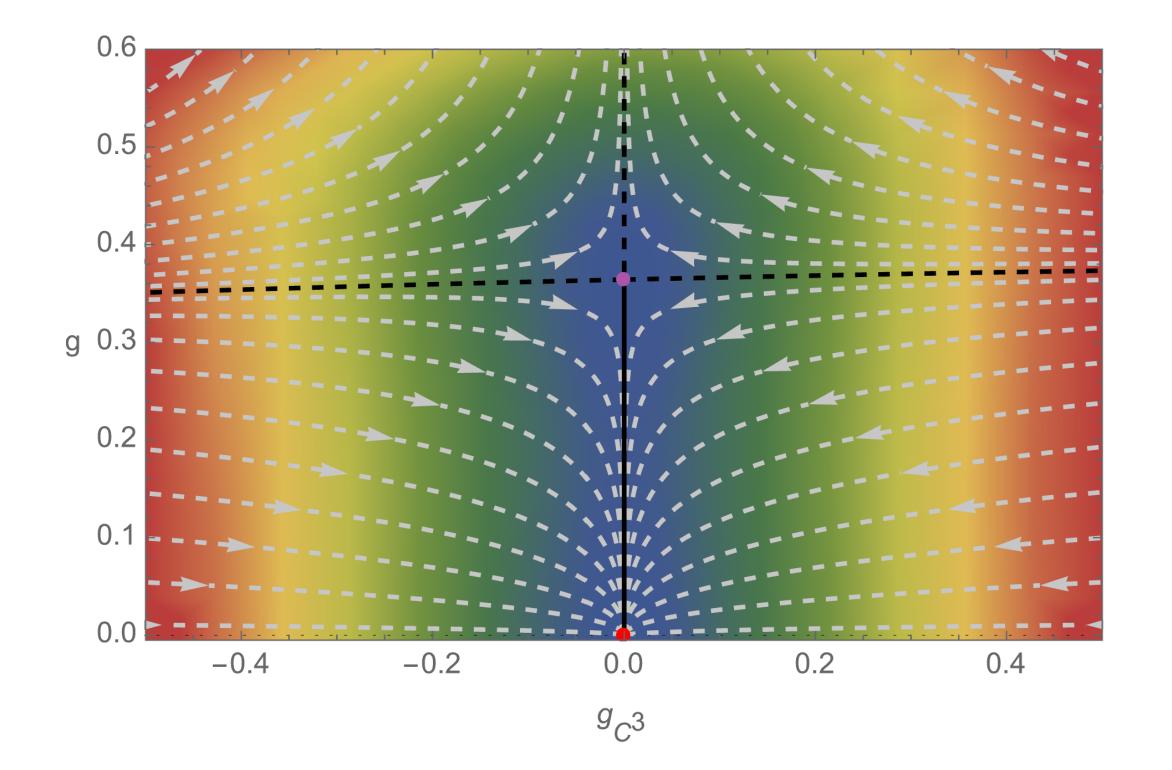
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10 running field redefinitions!

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$$\theta_1 = 2.225$$
 $\theta_2 = -3.850$

consider derivative expansion up to 6th order (d=4):

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$$G_{C^3} = \left(\frac{G_N}{16\pi}\right)^2 \left(a - \frac{86}{315} \ln G_N k^2\right)$$
$$a = -0.5065$$

- work in progress:
 - curvature expansion up to second order (two momentum-dependent running field redefinitions)

 BK, Ripken
 - "f(R)" in [GR]

 BK, Sannestedt
 - general dimension, gauge dependence, on-shell flow

consider derivative expansion up to 4th order coupled to matter:

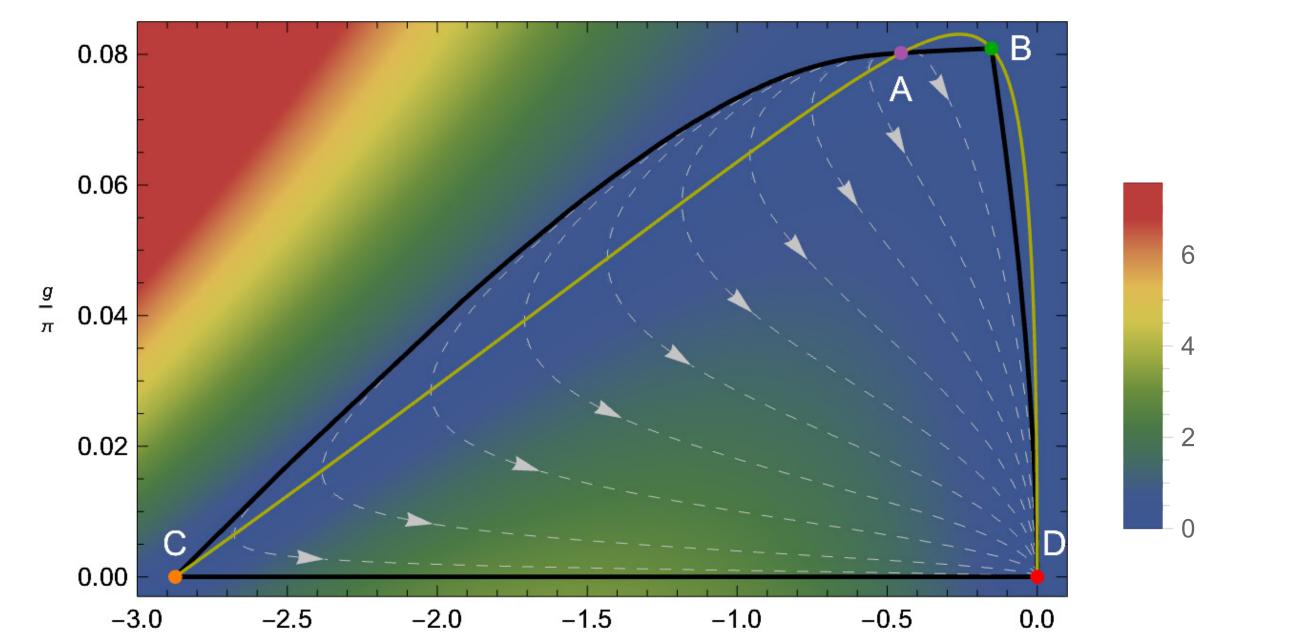
$$\Gamma_{[\mathbf{GR}]+[\phi]} = \int d^4x \left(\frac{1}{16\pi G_N} \left[2\Lambda - R + \Theta \mathfrak{E} \right] + \frac{1}{2} (D_\mu \phi)^2 + G_{D\phi^4} \left(\frac{1}{2} (D_\mu \phi)^2 \right)^2 \right)$$

 BK '22

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BK '22



A+C potentially fake, B+D stable [de Brito, BK, Schiffer '23]

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BK '22

$$\Gamma_{[\mathbf{GR}]+[\gamma]} = \int d^4x \left(\frac{1}{16\pi G_N} \left[2\Lambda - R + \Theta \mathfrak{E} \right] + \frac{1}{4} \text{tr} F^2 + G_{\mathcal{F}^2} \left(\frac{1}{4} \text{tr} F^2 \right)^2 + \frac{G_{F^4}}{16} \text{tr} F^4 + G_{CFF} C^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma} \right)$$

BK, Platania (WIP)

very similar phase diagram potentially intriguing relation to positivity/weak-gravity bounds

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BK '22

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BK, Platania (WIP)

many more systems under investigation stay tuned!

Summary

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- falsifiability is at the heart of science, and it should also be at the heart of quantum gravity research
- scattering amplitudes are promising tool to probe quantum gravity
- field redefinitions allowed all the difficulty starts at four-point function, strong results need high level of sophistication
- **[GR]** promising from perspective of Asymptotic Safety maybe no free parameters? AS **[GR]** = string theorists' dreams come true?